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MAGNETIC HELICITY IN THE SOLAR ATMOSPHERE: MUCH GAINED, STILL A LOT TO LEARN

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RCAAM OF THE ACADEMY OF ATHENS



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From Pioneering Thinking and Reasoning …

- Gauge-invariance theorem by Woltjer (1958)
- $\int_{V} \mathbf{A} \cdot \operatorname{curl} \mathbf{A} dV = \operatorname{constant},$
- Taylor's min-energy works in the seventies (1974, 1976)
- Mitch's decompositions / flux calculations in the eighties (1984)
- N. Seehafer's current helicity in1990
- Calugareanu invariant of Moffatt & Ricca (1992)
- Assertion on CMEs by B. C. Low (1994)
- Alexei's major observational finding in 1995
- D. Rust's and A. Kumar's works in 1996

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D. Canfield's and co. work on sigmoids in 1999

Many researchers in the 2000s and 2010s, aiming toward a practical calculation: H. Zhang, K. Kusano, P. Demoulin, A. Nindos, B. J. LaBonte, M. Zhang, K. Kuzanyan, E. Pariat, G. Valori, S. Regnier, J. Thalmann, A. Yeates, Y. Guo ...

... to Recent Developments

Formation of helical magnetic flux ropes prior to eruptions



Courtesy: George Chintzoglou (LMSAL)

Also, Chintzoglou et al., (2015); Nindos et al., (2015)

(see also works on helical flux ropes by S. Gibson)



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 Conservation of magnetic helicity in CMEs and extrapolation to get the CME Bz at L1



Patsourakos & Georgoulis (2016)

(see also works on helical flux ropes by S. Gibson)



Instantaneous Magnetic Helicity Budget vs. Helicity Injection Rate

 Knowledge of the 3D field above a boundary allows inference of the helicity budget

$$H = \int_V \mathbf{A} \cdot \mathbf{B} dV$$

gaugedependent in general

 Subtracting the reference helicity from the potential field, allows calculation of the relative magnetic helicity

$$H = \int_{V} (\mathbf{A} + \mathbf{A}_{\mathbf{p}}) \cdot (\mathbf{B} - \mathbf{B}_{\mathbf{p}}) dV$$

and alternative approaches, such as the field-line helicity (e.g., Aly, FDR, Lowder & Yeates, 2017)

$$\mathcal{A}(L) = \int_{L(x)} \frac{\mathbf{A} \cdot \mathbf{B}}{|\mathbf{B}|} d\ell$$



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$$\mathcal{A}(L) = \int_{L(x)}^{\prime} \frac{\mathbf{A} \cdot \mathbf{B}}{|\mathbf{B}|} d\ell$$

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 On the other hand, knowledge of the velocity field and the magnetic field vectors on the boundary plane allows evaluation of the Poynting theorem for relative magnetic helicity (e.g., Berger & Field, 1984; Kusano et al., 2002)

$$\frac{dH}{dt} = 2 \int_{S} \mathbf{A} \times (\mathbf{u} \times \mathbf{B}) \cdot \hat{\eta} dS$$

and its practical implementation by Demoulin & Berger (2003):

$$\frac{dH}{dt} = -2 \int_{S} (\mathbf{A}_{\mathbf{p}} \cdot \mathbf{u}_{\mathbf{ct}}) B_{n} dS$$

 Then the relative helicity is obtained by time integration of (dH / dt)

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However, there are Caveats and Shortcomings in Both Approaches

 The unmeasured coronal field is ambiguous and non-unique from 3D field extrapolations, thus having an unknown effect on helicity







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However, there are Caveats and Shortcomings in Both Approaches

while the velocity field vector on the boundary is also unknown and ambiguous, plus helicity injection rate calculations lack a point of reference



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Helicity Between Global and Local Solar Scales

 Solar magnetic helicity is a global quantity, but is mostly contributed by local (i.e., active region) scales
 Georgoulis et al., (2009)



- 80% of helicity stems from peculiar active-region flows; the rest from solar differential rotation
- > 99% of helicity stems from active regions; the rest from the quiet Sun

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How Can We Reconcile Between Scales, Uncertainties, Caveats and Shortcomings?

- Treat (relative) magnetic helicity <u>self-consistently</u> with (free) magnetic energy
- If possible, disentangle detailed knowledge of the 3D field from calculation
- Define test cases assess similarities and differences between methods



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LFF field approach

$$E = \mathcal{R}_{\ell} \lambda^2 \alpha^2 E_p$$

 $H = 8\pi \mathcal{R}_{\ell} \lambda^2 \alpha E_p$

$$H = \frac{8\pi}{\alpha}E$$

 $\alpha = const.$; $\lambda \longrightarrow length$ element

$$\mathcal{R}_{\ell} = \frac{1}{2} \sum_{l=1}^{n_x} \frac{\sum_{m=1}^{n_y} |b_{u_l;v_m}^2| / (u_l^2 + v_m^2)^{3/2}}{\sum_{m=1}^{n_y} |b_{u_l;v_m}^2| / (u_l^2 + v_m^2)^{3/2}}$$

Georgoulis & LaBonte (2007)



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Georgoulis & LaBonte (2007)

NLFF field approach

$$E = \lambda^2 A \sum_{l=1}^{N} \alpha_l^2 \Phi_l^{2\delta} + \frac{1}{8\pi} \sum_{l=1}^{N} \sum_{m=1_{l \neq m}}^{N} \alpha_l \mathcal{L}_{lm}^{arch} \Phi_l \Phi_m$$
$$H = 8\pi \lambda^2 A \sum_{l=1}^{N} \alpha_l \Phi_l^{2\delta} + \sum_{l=1}^{N} \sum_{m=1_{l \neq m}}^{N} \mathcal{L}_{lm}^{arch} \Phi_l \Phi_m$$

- N flux tubes ; $\lambda \rightarrow$ length element; A, δ const.
- E is a <u>lower limit free energy</u> for a given connectivity that ignores intertwining of flux tubes in the corona (based on the analysis of Demoulin et al., (2006)

Georgoulis et al., (2012)

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Assessment of a Skeleton Connectivity, <u>Without</u> a Detailed 3D Knowledge of the Magnetic Field



- The normal photospheric field component is partitioned; each partition assumed a different flux tube
- Connectivity inferred via a simulated annealing scheme favoring shortest connections, i.e., alongside polarity inversion lines

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Can This Approach be Applied to Global Scales?

 It could, conceivably, utilizing spherical geometry on a synoptic <u>vector</u> magnetogram



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 This scheme would find the connections within active regions first, before connecting largest scales



Can This Approach be Applied to Global Scales?

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 This scheme would find the connections within active regions first, before connecting largest scales

This remains to be implemented

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Testing Helicity Calculation Methods on Synthetic Test Cases

ISSI Bern team on magnetic helicity and applications (G. Valori & E. Pariat, Team Leaders)

Finite volume (FV)	Helicity-flux integration (FI)						
$\begin{aligned} \mathcal{H}_{\mathcal{V}} &= \int_{\mathcal{V}} \left(\mathbf{A} + \mathbf{A}_{p} \right) \cdot \left(\mathbf{B} - \mathbf{B}_{p} \right) \mathrm{d}\mathcal{V} \\ &\text{see Eq. (3)} \end{aligned}$	$\frac{\mathrm{d}\mathcal{H}_{\mathcal{V}}}{\mathrm{d}t} = 2\int_{\partial\mathcal{V}} [(\mathbf{A}_{\mathrm{p}} \cdot \mathbf{B})v_n - (\mathbf{A}_{\mathrm{p}} \cdot \mathbf{v}_t)B_n]\mathrm{d}S$						
 Requires B in V e.g., from MHD simulations or NLFFF Compute H_V at one time May employ different gauges (see Table 2) 	 Requires time evolution of vector field on ∂V Requires knowledge or model of flows on ∂V Valid for a specific set of gauge and assumptions see Pariat et al, (2017) 						
Discrete flux-tubes (DT)							
$\mathcal{H} \simeq \sum_{i=1}^{M} \mathcal{T}_i \Phi_i^2 + \sum_{i=1}^{M} \sum_{j=1, j \neq i}^{M} \mathcal{L}_{i,j} \Phi_i \Phi_j,$							
see Eq. (31)							
Twist-number (TN)	Connectivity-based (CB)						
${\cal H}{\simeq}{\cal T}{\Phi}^2$	$\mathcal{H} = A \sum_{i=1}^{M} \alpha_i \Phi_i^{2\delta} + \sum_{l=1}^{M} \mathcal{L}_{lm} \Phi_l \Phi_m$						
see Eq. (32)	see Eq. (35)						
 Estimation of the twist contribution to <i>H</i> Requires B in <i>V</i> Requires a flux-rope-like structure for computing – Models the corona connectivity as a collection of <i>M</i> force-free flux tubes Minimal connection length principle 							

Valori et al., (2016)



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Testing Helicity Calculation Methods on Synthetic Test Cases

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$\mathcal{H}_{\mathcal{V}} = \int_{\mathcal{V}} (\mathbf{A} + \mathbf{A}_{p}) \cdot (\mathbf{B} - \mathbf{B}_{p}) d\mathcal{V}$ see Eq. (3)	Method		Label	el Category		Section	Reference
 Requires B in V e.g., from MHD simulati NLFFF Compute H_V at one time May employ different gauges (see Table 2) 	Cou Tha	ılomb- lmann	Coulomb_JT	Finite v	olume	Sect. <u>2.1.1</u>	Thalmann et al. (<u>2011</u>)
$\begin{array}{l} Disc\\ \mathcal{H} \simeq \sum_{i=1}^{M} \mathcal{T}_{i} \Phi_{i}^{2}\\ Twist-number\ (TN)\\ \mathcal{H} \simeq \mathcal{T} \Phi^{2}\\ \text{see Eq.}\ (32)\\ \end{array}$ $\begin{array}{l} - \text{ Estimation of the twist contribution to } \mathcal{H}\\ - \text{ Requires } \mathbf{B} \text{ in } \mathcal{V}\\ - \text{ Requires a flux-rope-like structure for cor}\\ \text{the twist } \mathcal{T}\end{array}$	Coulomb-Yang		Coulomb_SY	Finite volume		Sect. 2.1.2	Yang et al. (<u>2013b</u>)
	Cou Ruc	ılomb- lenko	Coulomb_GR	Finite v	olume	Sect. <u>2.1.3</u>	Rudenko and Anfinogentov (<u>2014</u>)
	De	/ore-Valori	DeVore_GV	Finite v	olume	Sect. 2.2.1	Valori et al. (<u>2012</u>)
	DeVore-Moraitis		DeVore_KM	Finite v	olume	Sect. 2.2.2	Moraitis et al. (<u>2014</u>)
	De Anf	/ore- inogentov	DeVore_SA	Finite volume		Sect. <u>2.2.3</u>	Not available
Valori et al., (2016)	Twi	st-number	TN	Discrete tubes	e flux-	Sect. <u>2.3.1</u>	Guo et al. (<u>2010</u>)
	Cor bas	nectivity- ed	СВ	Discrete tubes	e flux-	Sect. <u>2.3.2</u>	Georgoulis et al. (2012)



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Comparison Results ^[1]



- Four selected cases
 - (a) Low & Lou FF equilibrium
 - (b) Titov & Demoulin FF equilibrium
 - (c) MHD stable model of Leake et al., (2013)
 - (d) MHD unstable model of Leake et al., (2013)

Valori et al., (2016)

Comparison of 3D finite-volume methods in the two MHD configurations



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0.3

0.2

0.1

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Comparison of 3D finite-volume methods in the two MHD configurations

Pretty good agreement between methods, if 3D field is known. Also, reasonably immune results to magnetic reconnection





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Comparison Results ^[2]



Less than an agreement in MHD stable configuration, but agreement within 10% for the MHD unstable one!

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 Comparison between the full 3D method and the one with skeleton connectivity







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Future Steps in Correlating Magnetic Helicity Calculation Methods

- The paper of Valori et al., (2016) is only Paper I
- In Paper II (Pariat et al., 2018, in prep.) helicity-injection rates will be tested
- In Paper III (Georgoulis et al., 2018, in prep.) methods will be tested on an NLFFextrapolated, observed active-region case
- In another published work (Guo et al., 2017) the twist-number helicity method is applied to a number of models (Titov & Demoulin, MHD models, etc.)



Issues to Work Out. I. the "Energy — Helicity" Diagram





 Active regions, eruptive or not, exhibit a distinctive scaling relation between free magnetic energy and absolute value of relative magnetic helicity

$H \propto E^{0.84 \pm 0.05}$

 This was later noticed for quiet-Sun structures, and even for MHD models, (Tziotziou et al., 2014)

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 Notice the jump to lower helicities in case of quiet-Sun structures — this points to an overall incoherence of helical sense in the quiet Sun, that might be expected, but need to be investigated further

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Issues to Work Out. II. Competition of the Two Helicity Senses

 Our discrete relative helicity calculation method enables the calculation of both signs of helicity within a given magnetic structure



 Active regions statistically show a <u>dominant</u> helicity sign, contrary to quiet-Sun regions



Issues to Work Out. II. Competition of the Two Helicity Senses

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 Both X-class flares relate to eruptions of righthanded structures (e.g., Jiang et al., 2013)

- NOAA AR 11283, on Sep 2011
- Initially left-handed configuration
- Gives two eruptive X-class flares with very low helicity
- Helicity annihilation (Kusano et al., 2003)?



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- Both X-class flares relate to eruptions of righthanded structures (e.g., Jiang et al., 2013)
- What seems to happen? Initially left-handed structure gradually turns into a right-handed one

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Issues to Work Out. III. Mutual vs. Self Helicity

 Our discrete relative helicity calculation method also enables separation between self and mutual terms of relative helicity and free energy

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Issues to Work Out. III. Mutual vs. Self Helicity

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Other Helicity Realizations: the "Current-Carrying" Helicity

• The "current-carrying" helicity of Pariat et al., (2017)

$$H_{V} = H_{j} + 2H_{pj} \quad \text{with}$$

$$H_{j} = \int_{\mathcal{V}} (A - A_{p}) \cdot (B - B_{p}) \, d\mathcal{V}$$

$$H_{pj} = \int_{\mathcal{V}} A_{p} \cdot (B - B_{p}) \, d\mathcal{V},$$
(10) "Current-carrying" helicity
$$H_{pj} = \int_{\mathcal{V}} A_{p} \cdot (B - B_{p}) \, d\mathcal{V},$$
(11) Mutual helicity



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Other Helicity Realizations: the "Current-Carrying" Helicity

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- (9) Relative helicity
- (10) "Current-carrying" helicity

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(11) Mutual helicity

The ratio $|H_j / H_V|$ seems to spike prior to the eruption in the simulation of Leake et al. (2013), implying a possible physical role for H_j

From its construction and the DeVore gauge (A.n =0), H_j does not have a contribution on ∂V and is scale invariant

Other Helicity Realizations: the "Current-Carrying" Helicity

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> More work is needed to understand H_j better

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Finally, the Helicity Spectra



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- NOAA AR 11158
- Realizability condition: $k|H(k,t)| \leq 2E(k,t)$

(constraining magnetic helicity, albeit in Fourier space)

 First complete calculation of the current helicity, albeit in Fourier space

$$H_c(k,t) \simeq k^2 H(k,t)$$

Hoping to hear more in this Thinkshop!

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Conclusions

- Ground-breaking progress in solar magnetic helicity studies over recent decades
- Relative] magnetic helicity is now placed on equal footing with [free] magnetic energy in solar low-atmospheric configurations
- As in every such progress, however, more questions than answers are borne. In particular:





Conclusions

- Ground-breaking progress in solar magnetic helicity studies over recent decades
- Relative] magnetic helicity is now placed on equal footing with [free] magnetic energy in solar low-atmospheric configurations
- As in every such progress, however, more questions than answers are borne. In particular:
 - We need to make sense / correlate between different helicity "flavors": relative (magnetic), "current-carrying", "spectral", current, kinetic, quadratic, cross-helicity, etc.
 - We need to understand the interplay between the two different senses of magnetic helicity in the <u>same</u> magnetic structure.
 - We need to understand the interplay between the mutual and selfhelicity in emerging magnetic structures.
 - Coherence of active-region helicity (deep-seated) vs. randomness of QS helical patterns (near-surface[?])

Thank you!



